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Complex dimensions and lengths of epsilon-neighborhoods of orbits

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Talk based on

P. Mardešić, G. Radunovic, M. Resman, *Fractal zeta functions of orbits of parabolic diffeomorphisms*, submitted (2020), https://arxiv.org/abs/2010.05955

Standard zeta function

• The standard Hurwitz (Riemann) zeta function

$$\zeta_a(s) := \sum_{j=0}^{\infty} \frac{1}{(j+a)^s}, \ a > 0, \ \mathsf{Re}(s) > 1$$

- ullet converges absolutely for $\mathrm{Re}(s)>1$
- meromorphically extendable to $\mathbb{C} \setminus \{1\}$
- single pole at 1 with residue $\mathrm{Res}(\zeta_a(s),s=1)=1$
- for a=1: the Riemann zeta function

'Geometric generalizations' - **fractal zeta functions** in the sense of *Lapidus*

- $\mathcal{L}:=\{\ell_j:\,j\in\mathbb{N}\}$ a disjoint union of intervals on the real line with lengths ℓ_j
- (1) The geometric zeta function of a fractal string (Lapidus, Frankenhuijsen, 2000)

$$\zeta_{\mathcal{L}}(s):=\sum_{j=1}^{\infty}\ell_{j}^{s},\ s\in\mathbb{C}$$
 s.t. the sum converges absolutely

 $\star \ \ell_j := rac{1}{j}$ standard zeta function



Generalizations for arbitrary sets

(2) The distance zeta function of a bounded set $A \subseteq \mathbb{R}^N$

$$\zeta_A(s) := \int_{A_\delta} d(x, A)^{s-N} dx$$

- $\delta > 0$ inessential (up to a holomorphic function)
- (3) The tube zeta function of a bounded set $A \subseteq \mathbb{R}^N$:
 - the tube function of *A*:

$$\varepsilon\mapsto V_A(\varepsilon):=|A_\varepsilon|$$
 (the Lebesgue measure)

• $V_A(\varepsilon) \sim M \varepsilon^{N-s_0}, \ \varepsilon \to 0 \Rightarrow \dim_B(A) = s_0, \ \mathcal{M}^{s_0}(A) = M.$

$$\tilde{\zeta}_A(s) := \int_0^\delta t^{s-N-1} V_A(t) \, dt,$$

$$\operatorname{Re}(s) > \dim_B(A), \ \delta > 0$$
 inessential

(Lapidus, Frankenhuijsen 2000, 2006; Lapidus, Radunović, Žubrinić, 2017)

For fractal strings, all three equal up to a holomorphic function

$$\mathcal{L} \Rightarrow A := \{a_j : j \in \mathbb{N}_0\}, \ \ell_j := a_{j-1} - a_j$$

The functional equations on domains of definition (up to holomorphic functions):

- $\zeta_A(s) = \frac{2^{N-s}}{s} \zeta_{\mathcal{L}}(s)$,
- $\tilde{\zeta}_A(s) = \frac{1}{N-s} \zeta_A(s)$, $\operatorname{Re}(s) > \dim_B(A)$.

Definition

Let

- $A \subseteq \mathbb{R}^N$ bounded,
- $\zeta_A(s)$ admits the meromorphic extension to whole $\mathbb C.$

The set of all poles is called the set of complex dimensions of A, $\Omega(A)$.

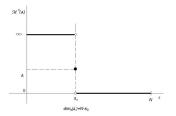
- $\zeta_A(s)$ holomorphic for $\operatorname{Re}(s) > \dim_B(A)$,
- simple pole at $s = \dim_B(A)$.

Complex dimensions (and their residues i.e. principal parts) *talk* about the geometry of the set! Similarly as the tube function!

The **box dimension** of a set (type of fractal dimension)

- $A \subseteq \mathbb{R}^N$ bounded
- $\varepsilon > 0$, $V_A(\varepsilon) := |A_{\varepsilon}|$ the Lebesque measure of the ε -neighborhood (the so-called *tube-function*)
- For $s \in [0, N]$, consider (or: \liminf, \limsup)

$$\lim_{\varepsilon \to 0} \frac{V_A(\varepsilon)}{\varepsilon^{N-s}} \in [0, \infty],$$



Slika:
$$s\mapsto \lim_{\varepsilon\to 0} \frac{V_A(\varepsilon)}{\varepsilon^{N-s}}$$
, $s\in [0,N]$.



- the moment of jump $s_0 \equiv$ the box dimension, $\dim_B(A) = s_0$.
- the value at $s_0 \equiv \text{Minkowski content}$, $\mathcal{M}(A)$.
- if $V_A(\varepsilon) \sim_{\varepsilon \to 0} C\varepsilon^{N-s_0} \Rightarrow \dim_B(A) = s_0, \ \mathcal{M}(A) = C.$

Example 1

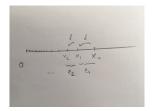
- $\dim_B(point) = 2 2 = 0;$
- $\dim_B(finite length\ line) = 2 1 = 0;$
- $\dim_B(A) = 2 0 = 2, |A| > 0.$

Non-trivial examples: Orbits of local parabolic diffeomorphisms (\equiv germs) on the real line \mathbb{R}_+

- (attracting) parabolic germ $f(x) = x ax^{k+1} + \ldots \in \text{Diff}(\mathbb{R}_+, 0), \ a > 0, \ k \in \mathbb{N}$ $a_j \sim j^{-1/k}, \ j \to \infty,$
- (attracting) hyperbolic germ $f(x) = \lambda x + ..., 0 < \lambda < 1$ $a_j \sim \lambda^j, \ j \to \infty.$

Orbit of f with initial point $x_0 \in (\mathbb{R}_+, 0)$:

$$\mathcal{O}_f(x_0) := \{ x_n := f^{\circ n}(x_0) : n \in \mathbb{N}_0 \}$$





Box dimension and Minkowski content of orbits

Žubrinić, Županović 2005, MRŽ 2012

ullet a parabolic orbit of *multiplicity* k

$$V_{\mathcal{O}^f(x_0)}(\varepsilon) \sim (2/a)^{\frac{1}{k+1}} \frac{k+1}{k} \varepsilon^{\frac{1}{k+1}} + o(\varepsilon^{\frac{1}{k+1}}), \ \varepsilon \to 0,$$

$$\dim_B(\mathcal{O}^f(x_0)) = 1 - \frac{1}{k+1}, \ \mathcal{M}(\mathcal{O}^f(x_0)) = (2/a)^{\frac{1}{k+1}} \frac{k+1}{k},$$

a hyperbolic orbit

$$V_{\mathcal{O}^f(x_0)}(\varepsilon) \sim a(\lambda) \cdot \varepsilon(-\log \varepsilon) + o(\varepsilon(-\log \varepsilon)), \ \varepsilon \to 0,$$

$$\dim_B(\mathcal{O}^f(x_0)) = 1 - 1 = 0, \ \mathcal{M}(\mathcal{O}^f(x_0)) = +\infty.$$

Later: R [2013]

* formal class of f using asymptotic expansion of function

$$\varepsilon \mapsto V_{\mathcal{O}^f(x_0)}(\varepsilon)$$
, as $\varepsilon \to 0$

* further complex dimensions needed



Example 2 (The complex dimensions of the ternary Cantor set, LRŽ 2017)

 \star viewed as a fractal string, the order of intervals not important $\mathcal{L}_{\mathcal{C}}\text{, }A$

$$\zeta_{\mathcal{L}_{\mathcal{C}}}(s) = \sum_{j=1}^{\infty} \ell_{j}^{s} = \sum_{k=0}^{\infty} 2^{k} \left(\frac{1}{3^{k+1}}\right)^{s} = \frac{1}{3^{s}-2}, \ |\frac{2}{3^{s}}| < 1$$

- holomorphic for $\operatorname{Re}(s) > \log_2 3 = \dim_B \mathcal{C}$
- ullet unique meromorphic extension to ${\mathbb C}$ by the above formula with poles:

$$\Omega(\mathcal{C}) = \{ \omega_k := \log_3 2 + i \frac{2k\pi}{\log 3}, \ k \in \mathbb{Z} \}.$$

Example 3 (The tube function of the Cantor set (LRŽ 2017))

$$V_{\mathcal{C}}(\varepsilon) = \varepsilon^{1-\log_3 2} (G(-\log \varepsilon) + o(1)), \ \varepsilon \to 0,$$

G a nonconstant periodic function.

A conjecture (LRŽ): Strong oscillations in the first term indication of self-similarity; non-real complex dimensions; possible definition of *fractality* of a set as possessing non-real complex dimensions?

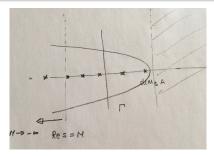
Complex dimensions deduced from asymptotics of the tube function of a set

(formally proven in LRŽ, 2017)

- $\star\ \tilde{\zeta}_A$ the tube zeta function of set $A\subseteq\mathbb{R}^N,$ meromorphically extendable to $\mathbb{C}.$
- $\star t \mapsto V_A(t) = |A_t|$, $t \in (0, \delta)$, the tube function of A
 - $\tilde{\zeta}_A(s) = \mathcal{M}(\chi_{(0,\delta)}V_A/\mathrm{id}^N)(s) = \int_0^\delta V_A(t)t^{s-1-N} dt$
 - Conversely,

$$V_A(t) = \frac{t^N}{2\pi i} \mathcal{I} \mathcal{M}(\tilde{\zeta}_A)(t) = \frac{1}{2\pi i} \int_{\Gamma_c} \tilde{\zeta}_A(s) t^{N-s} \, ds, \ t \in (0, d).$$

 Γ ... a vertical line at around $s=c, \ c>\dim_B A$



* Heuristically, the residue theorem 'gives' expansions of $t\mapsto V_A(t)$ from poles and residues of ζ_A :

 \star e.g. $\Omega_A = \{\omega_n, n \in \mathbb{N}\}$ only first-order poles

$$(**) V_A(t) = \frac{1}{2\pi i} \int_{\Gamma_c} \frac{t^{N-s}}{N-s+1} \tilde{\zeta}_A(s) ds =$$

$$= \frac{1}{2\pi i} \sum_{\omega \in \Omega_A, \operatorname{Re}(\omega) > -M} \frac{t^{N-\omega}}{N-\omega+1} \operatorname{Res}(\tilde{\zeta}_A, \omega) + O(t^{N+M}), t \to 0, M \in \mathbb{N}.$$

(in case of higher-order poles logarithmic terms in the expansion)

Idea of proof of (**) (LRŽ)

- to get asymptotic remainder $O(t^{N+M})$, $M\in\mathbb{N}$, bounds needed on zeta function along vertical lines $\mathrm{Re}(s)=-M,$ $M\to\infty$
- so-called *languidity bounds* of $\tilde{\zeta}_A(s)$ along vertical lines $s=\sigma+i au$, as $au \to \pm \infty$
- pointwise asymptotics as long as bounds rational

$$|\tilde{\zeta}_A(\sigma + i\tau)| \sim \tau^{-\gamma}, \ \gamma > 0, \tau \to \pm \infty$$

- polynomial bounds $(\gamma < 0) \Rightarrow$ only distributional asymptotics (there exists some primitive of tube function $t \mapsto V_A^{[k]}$) that expands pointwise up to this term, but differentiation of asymptotic expansions can be done just distributionally!
- $\frac{t^{N-s+k}}{(N-s+1)_k} \tilde{\zeta}_A(\sigma+i\tau)$, as $au o \pm \infty$, becomes rational for k sufficiently big!



Relation to dynamical systems - 1-D orbits of germs seen as fractal strings

• (attracting) parabolic germ $f(z) = z - ax^{k+1} + \ldots \in \text{Diff}(\mathbb{R}_+, 0), \ a > 0, \ k \in \mathbb{N}$

$$a_j \sim j^{-1/k}, \ \ell_j \sim j^{-\frac{k+1}{k}}, \ j \to \infty,$$

• (attracting) hyperbolic germ $f(x) = \lambda x + ..., 0 < \lambda < 1$

$$a_j \sim \lambda^j, \ \ell_j \sim \lambda^j, \ j \to \infty.$$

$$\zeta_{\mathcal{L}_f}(s)$$
" \sim " $\sum_{j \in \mathbb{N}} j^{-s\frac{k+1}{k}}$

- * holomorphic for $\operatorname{Re}(s) > \frac{k}{k+1} = \dim_B \mathcal{O}^f(x_0)$
- however, too coarse approximations for meromorphic extensions - info on poles and residues lost
- * notation: $\zeta_{\mathcal{L}_f},\,\zeta_f,\,\widetilde{\zeta}_f$



Precise computations tedious even in the simplest model case of germs, k=1, $\rho=0$ (MRR 2020)

- * Model cases with residual invariant ho=0 and multiplicity $k\in\mathbb{N}$
- * time-one maps of simple vector fields $x' = -x^{k+1}$:

$$f_k(x) := \operatorname{Exp}(x^{k+1} \frac{d}{dx}) = \frac{x}{(1+kx^k)^{1/k}} = x - x^{k+1} + o(x^{k+1}), \ k \in \mathbb{N}.$$

Putting $X := x_0^{-1}$,

$$\ell_j = \frac{1}{(j+X)(j+1+X)} = \frac{1}{(j+X)^2} \cdot \left(1 + \frac{1}{j+X}\right)^{-1},$$

$$\ell_j^s = \frac{1}{(j+X)^{2s}} \cdot \left(1 + \frac{1}{j+X}\right)^{-s} =$$

$$= \sum_{m=0}^{\infty} {\binom{-s}{m}} \frac{1}{(j+X)^{2s+m}}.$$

Heuristically (formal change of order of summation),

$$\zeta_{\mathcal{L}_{f_1}}(s) = \sum_{j=0}^{\infty} \ell_j^s \, \sim \sum_{m=0}^{\infty} {\binom{-s}{m}} \zeta_X(2s+m).$$
(1)

Complex dimensions: $\omega_n := \frac{1-n}{2}, \ n \in \mathbb{N}_0$, with residues:

$$\operatorname{Res}(\zeta_{\mathcal{L}_{f_1}},\omega_n)=inom{n-1}{2}.$$
 Zero residue for n odd.

What to do in the case $\rho \neq 0$ or even non-model case?

Arbitrary parabolic germ

$$f(x) = x - ax^{k+1} + o(x^{k+1}) \in \text{Diff}(\mathbb{R}_+, 0)$$

Theorem B (MRR 2020, Complex dimensions for arbitrary parabolic orbits)

 $f \in \mathrm{Diff}(R_+,0)$, of formal class (k,ρ) , $k \in \mathbb{N}$, $\rho \in \mathbb{R}$.

(1) The distance zeta function $\zeta_f(s)$ can be meromorphically extended to \mathbb{C} .

Theorem B

(2) For $s \in W_M := \{s > 1 - \frac{M}{k+1}\}, M \in \mathbb{N}$:

$$\zeta_f(s) = (1-s) \sum_{m=1}^k \frac{a_m}{s - \left(1 - \frac{m}{k+1}\right)} + (1-s) \left(\frac{b_{k+1}(x_0)}{s} + \frac{a_{k+1}}{s^2}\right) + \left(1-s\right) \sum_{m=k+2}^{M-1} \sum_{p=0}^{\lfloor \frac{m}{k} \rfloor + 1} \frac{(-1)^p p! \cdot c_{m,p}(x_0)}{\left(s - \left(1 - \frac{m}{k+1}\right)\right)^{p+1}} + g(s),$$

g(s) holomorphic in W_M .

- st the coefficients in principal parts of poles real, with dependence on x_0 , as noted!
- * related to the coefficients of the asymptotic expansion of the tube function of the orbit!
- * new wrt model: higher-order poles correspond to logarithmic terms in the asymptotic expansion of the tube function due to $\rho \neq 0$



Hyperbolic orbits

$$* \mathcal{O}_{f}(x_{0}) = \{x_{0}\lambda^{n} : n \in \mathbb{N}_{0}\},$$

$$* \mathcal{L}_{f} := \{\ell_{j} := f^{\circ j}(x_{0}) - f^{\circ (j+1)}(x_{0}) = x_{0}(1-\lambda)\lambda^{j} : j \in \mathbb{N}_{0}\},$$

$$*$$

$$\zeta_{f}(s) := \frac{2^{1-s}}{s} \sum_{j=0}^{\infty} \ell_{j}^{s} = \frac{2^{1-s}x_{0}^{s} \cdot (1-\lambda)^{s}}{s} \frac{1}{1-\lambda^{s}},$$

* extends meromorphically from $\{s\in\mathbb{C}: \operatorname{Re}(s)>0\}$ to \mathbb{C} : double pole $s_0=0$ and the simple poles

$$s_k := \frac{2k\pi}{\log \lambda} i, \ k \in \mathbb{Z}.$$

*

$$V_f(\varepsilon) = -\frac{2}{\log \lambda} \varepsilon (-\log \varepsilon) + H\left(\log_{\lambda} \frac{2\varepsilon}{x_0(1-\lambda)}\right) \cdot \varepsilon,$$

 $H:[0,+\infty)\to\mathbb{R}$ a 1-periodic bounded function

Generalized asymptotic expansion of the tube function of orbits $\varepsilon \mapsto V_f(\varepsilon)$

$$f(x)=x-ax^{k+1}+o(x^{k+1})\in \mathrm{Diff}(\mathbb{R},0),\ a>0,$$
 arbitrary parabolic germ

'Problem' noted in (R, 2014), (MRRZ, 2019):

- (*) the tube function $\varepsilon\mapsto V_f(\varepsilon)$ fails to have a full asymptotic expansion in power-logarithm scale,
- (*) oscillatory coefficient at order $O(\varepsilon^{\frac{2k+1}{k+1}}), \ \varepsilon \to 0.$

Generalized asymptotic expansion of the tube function of orbits $\varepsilon \mapsto V_f(\varepsilon)$

Proposition (MRR 2020)

A generalized asymptotic expansion of the tube function with full description of oscillatory coefficients:

$$\begin{split} V_f(\varepsilon) &\sim 2^{\frac{1}{k+1}} a^{-\frac{1}{k+1}} \frac{k+1}{k} \cdot \varepsilon^{\frac{1}{k+1}} + \sum_{m=2}^k a_m \cdot \varepsilon^{\frac{m}{k+1}} + 2\rho \frac{k-1}{k} \cdot \varepsilon \log \varepsilon + b_{k+1}(x_0) \varepsilon + \\ &+ \sum_{m=k+2}^{2k} \sum_{p=0}^{\lfloor \frac{m}{k} \rfloor + 1} c_{m,p} \varepsilon^{\frac{m}{k+1}} \log^p \varepsilon + \sum_{p=1}^{\lfloor \frac{2k+1}{k} \rfloor + 1} c_{2k+1,p} \varepsilon^{\frac{2k+1}{k+1}} \log^p \varepsilon + \\ &+ \tilde{P}_{2k+1}(G(\tau_\varepsilon)) \cdot \varepsilon^{\frac{2k+1}{k+1}} + \sum_{m=2k+2}^{\infty} \sum_{p=0}^{\lfloor \frac{m}{k} \rfloor + 1} \tilde{Q}_{m,p}(G(\tau_\varepsilon)) \cdot \varepsilon^{\frac{m}{k+1}} \log^p \varepsilon, \ \varepsilon \to 0^+. \end{split}$$

- (*) $\varepsilon \mapsto \tau_{\varepsilon}$ the so-called continuous critical time (MRRZ 2019)
- (*) $G:[0,+\infty)
 ightarrow \mathbb{R}$ 1-periodic, $G(s)=1-s, \ s\in (0,1), \ G(0)=0$
- (*) \tilde{P}_{2k+1} resp. $\tilde{Q}_{m,p}$, polynomials whose coefficients in general depend on coefficients of f and initial condition x_0 .

Two ways to 'regularize' oscillatory coefficients:

- the continuous-time tube function ('dynamical' regularization)
- the expansion of the tube function in the sense of Schwarz distributions

The continuous tube function-a 'dynamical smoothening' of the expansion (MRRZ 2019)

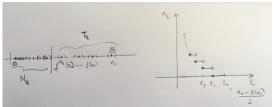
$$V_f(\varepsilon) = |O^f(x_0)_{\varepsilon}| = |T_{\varepsilon}| + |N_{\varepsilon}| = 2\varepsilon \cdot n_{\varepsilon} + (f^{n_{\varepsilon}}(x_0) + 2\varepsilon), \ \varepsilon > 0$$
 (the standard computation of the tube function: dividing into the *tail* and the *nucleus*)

(*) $\varepsilon \mapsto n_{\varepsilon}$ the so-called **discrete critical time** 'separating' tail and nucleus - jump function at $\varepsilon_n = \frac{f^n(x_0) - f^{n+1}(x_0)}{2} \to 0$, $n \to \infty$.

$$(*)$$
 $n_{\varepsilon}\in\mathbb{N}$ determined by two inequalities:

$$|f^{n_{\varepsilon}-1}(x_0) - f^{n_{\varepsilon}}(x_0)| \ge 2\varepsilon,$$

$$|f^{n_{\varepsilon}}(x_0) - f^{n_{\varepsilon}+1}(x_0)| < 2\varepsilon.$$





- (*) the continuous critical time (MRRZ 2019) $\varepsilon \mapsto \tau_{\varepsilon}$
- an analytic, dynamical 'approximation' of n_{ε}
- relies on embedding of f as the time-one map in a flow $\{f^t:\ t\in\mathbb{R}\}$:

$$f^{\tau_{\varepsilon}}(x_0) - f^{\tau_{\varepsilon}+1}(x_0) = 2\varepsilon.$$

Note: $n_{\varepsilon} = \lfloor \tau_{\varepsilon} \rfloor + 1$, $\varepsilon > 0$

More precisely, $n_{\varepsilon} = \tau_{\varepsilon} + G(\tau_{\varepsilon})$.

The continuous tube function $\varepsilon\mapsto V^c_f(\varepsilon)$

$$V_f^c(\varepsilon) = 2\varepsilon\tau_\varepsilon + (f^{\tau_\varepsilon}(x_0) + 2\varepsilon), \ \varepsilon > 0.$$

- $(*) \text{ analytic in } \varepsilon \in (0,\delta)$
- (*) expansion coincides with the expansion of $\varepsilon\mapsto V_f(\varepsilon)$ up to the first oscillatory term
- (*) full asymptotic expansion in a power-log scale, no oscillatory coefficients!

Asymptotic expansion of the continuous tube function

f parabolic

Proposition (MRR 2020)

$$V_f^c(\varepsilon) \sim 2^{\frac{1}{k+1}} a^{-\frac{1}{k+1}} \frac{k+1}{k} \cdot \varepsilon^{\frac{1}{k+1}} + \sum_{m=2}^k a_m \cdot \varepsilon^{\frac{m}{k+1}} + 2\rho \frac{k-1}{k} \cdot \varepsilon \log \varepsilon +$$

$$+ b_{k+1}(x_0) \varepsilon + \sum_{m=k+2}^{\infty} \sum_{p=0}^{\lfloor \frac{m}{k} \rfloor + 1} c_{m,p} \varepsilon^{\frac{m}{k+1}} \log^p \varepsilon, \ \varepsilon \to 0^+.$$

- (*) $c_{2k+1,0}$ resp. $c_{m,p}$, $m\geq 2k+2,\ p=0,\ldots,\lfloor \frac{m}{k}\rfloor+1$, are free coefficients of polynomials \tilde{P}_{2k+1} resp. $\tilde{Q}_{m,p}$
- (*) only the coefficient $b_{k+1}(x_0)$ depends on the initial condition x_0 .

Distributional expansion: distributional smoothening of the tube function

f parabolic

Proposition (MRR 2020)

$$V_f(\varepsilon) \sim_{\mathcal{D}} 2^{\frac{1}{k+1}} a^{-\frac{1}{k+1}} \frac{k+1}{k} \cdot \varepsilon^{\frac{1}{k+1}} + \sum_{m=2}^k a_m \cdot \varepsilon^{\frac{m}{k+1}} + 2\rho \frac{k-1}{k} \cdot \varepsilon \log \varepsilon + b_{k+1}(x_0)\varepsilon + \sum_{m=k+2}^{2k} \sum_{p=0}^{\lfloor \frac{m}{k} \rfloor + 1} c_{m,p} \varepsilon^{\frac{m}{k+1}} \log^p \varepsilon + \sum_{p=1}^{\lfloor \frac{2k+1}{k} \rfloor + 1} c_{2k+1,p} \varepsilon^{\frac{2k+1}{k+1}} \log^p \varepsilon + \sum_{m=2k+2}^{\lfloor \frac{m}{k} \rfloor + 1} c_{2k+1,p} \varepsilon^{\frac{2k+1}{k+1}} \log^p \varepsilon + \sum_{m=2k+2}^{\lfloor \frac{m}{k} \rfloor + 1} d_{m,p}(x_0) \cdot \varepsilon^{\frac{m}{k+1}} \log^p \varepsilon, \ \varepsilon \to 0^+.$$

Here.

$$\begin{array}{l} d_{2k+1,0}(x_0) := \int_0^1 \tilde{P}_{2k+1}(s) \, ds, \\ d_{m,p}(x_0) := \int_0^1 \tilde{Q}_{m,p}(s) \, ds, \ m \ge 2k+2, \ p = 0, \dots, \left\lfloor \frac{m}{k} \right\rfloor + 1, \end{array}$$

the mean values of 1-periodic functions $\tilde{P}_{2k+1} \circ G$ and $\tilde{Q}_{m,p} \circ G$.

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Thank you for your attention!